

Generalized measurements, quantum designs and 'quantum simplices'

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A generalized (discrete) quantum measurement:

POVM (positive operator valued measure) $\Pi = \{\Pi_j\}_{j=1}^k$, where

$$\Pi_j \in \mathcal{L}(\mathbb{C}^d), \quad \Pi_j \geq 0 \quad \text{and} \quad \sum_{j=1}^k \Pi_j = \mathbb{I}.$$

For pre-measurement state ρ , the **probability** of j -th outcome is $\text{tr}(\rho\Pi_j)$.

How to quantify the indeterminacy of the measurement?

The **Shannon entropy of POVM Π** is defined by:

$$H_{\Pi}(\rho) := \sum_{j=1}^k \eta(\text{tr}(\rho\Pi_j)),$$

for an initial state ρ , where $\eta(x) := -x \ln x$ ($x > 0$), $\eta(0) = 0$.

What are we looking for:

$$\min_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_{\Pi}(\tau) \quad \text{and} \quad \max_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_{\Pi}(\tau).$$

But that is the same as:

$$\min_{P \in \Delta_{\Pi}^1} H(P) \quad \text{and} \quad \max_{P \in \Delta_{\Pi}^1} H(P),$$

where $\Delta_{\Pi}^1 := \{P \in \Delta_k : P = (\text{tr}(\rho\Pi_1), \dots, \text{tr}(\rho\Pi_k)) \text{ for some } \rho \in \mathcal{P}(\mathbb{C}^d)\}$.

Let

$$\Delta_{\Pi} := \{P \in \Delta_k : P = (\text{tr}(\rho\Pi_1), \dots, \text{tr}(\rho\Pi_k)) \text{ for some } \rho \in \mathcal{S}(\mathbb{C}^d)\} = \text{conv}(\Delta_{\Pi}^1).$$

We call it the set of **‘allowed’ probabilities**.

How do Δ_{Π} and Δ_{Π}^1 look like?

How do Δ_Π and Δ_Π^1 look like?

- $\Delta_\Pi \subset \Delta_k$, $\dim \Delta_\Pi \leq d^2 - 1$ with equality iff Π is **informationally complete**
- Simplest case: Π is a rank-one PVM (projective valued measure). Then

$$\Delta_\Pi = \Delta_\Pi^1 = \Delta_k = \Delta_d.$$

- In general: **joint algebraic numerical range** of (Π_1, \dots, Π_k)
- In dimension 2: Δ_Π^1 is one of the following: a sphere, an ellipsoid, a disc, an ellipse (with the interior), a segment or a point.

When Δ_Π is 'the same' as $\mathcal{S}(\mathbb{C}^d)$?

For $A, B \in \mathcal{L}_s(\mathbb{C}^d)$ we denote their Hilbert-Schmidt inner product by $\langle\langle A|B \rangle\rangle$.

We say that Δ_Π is isomorphic to quantum state space $\mathcal{S}(\mathbb{C}^d)$ if there exists $\alpha > 0$ such that

$$\langle\langle \rho - I/d | \sigma - I/d \rangle\rangle = \alpha \langle p(\rho) - c, p(\sigma) - c \rangle \quad \forall \rho, \sigma \in \mathcal{S}(\mathbb{C}^d), \quad (1)$$

where $p(\rho) = (\text{tr}(\rho\Pi_1), \dots, \text{tr}(\rho\Pi_k))$ and $c = p(I/d) = (\text{tr}\Pi_1/d, \dots, \text{tr}\Pi_k/d)$.

Theorem

Δ_Π is isomorphic to quantum state space $\mathcal{S}(\mathbb{C}^d)$ iff $\{\Pi_i - (\text{tr}\Pi_i/d)I\}_{i=1}^k$ is a **tight frame** in $\mathcal{L}_s^0(\mathbb{C}^d)$. Then $\alpha = (d^2 - 1) / (\sum_{i=1}^k (\text{tr}\Pi_i^2 - (\text{tr}\Pi_i)^2/d))$.

V – N -dimensional Hilbert space with inner product $\langle \cdot | \cdot \rangle$

$F := \{f_1, \dots, f_m\} \subset V$

F is a **frame** if there exist $\alpha, \beta > 0$ such that

$$\alpha \|v\|^2 \leq \sum_{i=1}^m |\langle v | f_i \rangle|^2 \leq \beta \|v\|^2 \quad \text{for all } v \in V.$$

F is called a **tight frame** if $\alpha = \beta$.

An operator $S := \sum_{i=1}^m |f_i\rangle\langle f_i|$ is called a **frame operator**.

F is a tight frame iff $S = \alpha I$.

$\mathcal{L}_s(\mathbb{C}^d)$ – d^2 -dimensional real Hilbert space of selfadjoint operators on \mathbb{C}^d with the Hilbert-Schmidt inner product $\langle\langle A|B \rangle\rangle = \text{tr}(AB)$

$\mathcal{L}_s^0(\mathbb{C}^d) := \{A \in \mathcal{L}_s(\mathbb{C}^d) | \text{tr}A = 0\}$ – $(d^2 - 1)$ -dimensional subspace of traceless operators

$$\pi_0 : \mathcal{L}_s(\mathbb{C}^d) \ni A \mapsto A - (\text{tr}A/d)I \in \mathcal{L}_s(\mathbb{C}^d)$$

π_0 is the orthogonal projection onto $\mathcal{L}_s^0(\mathbb{C}^d)$

Theorem

Δ_Π is isomorphic to quantum state space $S(\mathbb{C}^d)$ iff $\{\Pi_i - (\text{tr}\Pi_i/d)I\}_{i=1}^k$ is a **tight frame** in $\mathcal{L}_s^0(\mathbb{C}^d)$. Then $\alpha = (d^2 - 1)/(\sum_{i=1}^k (\text{tr}\Pi_i^2 - (\text{tr}\Pi_i)^2/d))$.

Theorem (reformulated)

The following conditions are equivalent:

- Δ_Π is isomorphic to quantum state space $S(\mathbb{C}^d)$
- $\pi_0(\Pi)$ is a **tight frame** in $\mathcal{L}_s^0(\mathbb{C}^d)$ with frame constant $1/\alpha$
- $\sum_{i=1}^k |\pi_0(\Pi_i)\rangle\rangle\langle\langle\pi_0(\Pi_i)| = \frac{1}{\alpha}\mathcal{I}_0$.

Moreover, $\alpha = (d^2 - 1)/(\sum_{i=1}^k (\text{tr}\Pi_i^2 - (\text{tr}\Pi_i)^2/d))$.

Some special cases of POVMs Π such that $\pi_0(\Pi)$ is a tight frame in $\mathcal{L}_s^0(\mathbb{C}^d)$:

- $\text{tr}\Pi_i = d/k$ for $i = 1, \dots, k$. Then

$$\sum_{i=1}^k |\Pi_i\rangle\rangle\langle\langle\Pi_i| = \frac{1}{\alpha}\mathcal{I} + \frac{\alpha d - k}{\alpha dk} |I\rangle\rangle\langle\langle I|$$

In particular, Π is a **tight IC-POVM** [Scott '06]

- $\text{tr}\Pi_i = d/k$ **and** $\text{rank}\Pi_i = 1$ for $i = 1, \dots, k$. Then Π is a **2-design POVM**, i.e. pure states $\rho_i := \frac{k}{d}\Pi_i$ ($i = 1, \dots, k$) are such that

$$\frac{1}{k^2} \sum_{j,l=1}^k f(\text{tr}(\rho_j \rho_l)) = \iint_{\mathcal{P}(\mathbb{C}^d) \times \mathcal{P}(\mathbb{C}^d)} f(\text{tr}(\rho \sigma)) d\mu(\rho) d\mu(\sigma)$$

for every $f : \mathbb{R} \rightarrow \mathbb{R}$ polynomial of degree t or less, where μ denotes the unique unitarily invariant measure on $\mathcal{P}(\mathbb{C}^d)$.

Examples: SIC-POVMs, complete MUBs

- $\{\Pi_j\}_{j=1}^k$ is a 2-design POVM iff

$$\tau = (d+1) \sum_{j=1}^k \frac{k}{d} \text{tr}(\tau \Pi_j) \Pi_j - I$$

for every τ such that $\tau^* = \tau$ and $\text{tr} \tau = 1$.

- For 2-design POVM Π , $\rho_j = \frac{k}{d} \Pi_j$ and $p_j(\tau) = \text{tr}(\tau \Pi_j)$ we thus have

$$\sum_{j=1}^k p_j(\tau) = 1 \tag{2}$$

$$\sum_{j=1}^k (p_j(\tau))^2 = \frac{d(\text{tr}(\tau^2) + 1)}{k(d+1)} \tag{3}$$

$$\sum_{j,l,m=1}^k p_j(\tau) p_l(\tau) p_m(\tau) \text{tr}(\rho_j \rho_l \rho_m) = \frac{\text{tr}(\tau + I)^3}{(d+1)^3} \tag{4}$$

Remark: A self-adjoint operator τ is a pure quantum state iff $\text{tr}\tau = 1$, $\text{tr}\tau^2 = 1$ and $\text{tr}\tau^3 = 1$.

Theorem

Let Π be a 2-design POVM and let $\rho_j = \frac{k}{d}\Pi_j$ ($j = 1, \dots, k$). Let $(p_1, \dots, p_k) \in \mathbb{R}^k$ be such that

- $\frac{k}{d}p_l = (d+1) \sum_{j=1}^k p_j \text{tr}(\rho_j \rho_l) - 1$ for $l = 1, \dots, k$,
- $\sum_{j=1}^k p_j^2 = \frac{2d}{k(d+1)}$,
- $\sum_{j,l,m=1}^k p_j p_l p_m \text{tr}(\rho_j \rho_l \rho_m) = \frac{d+7}{(d+1)^3}$.

Then $(p_1, \dots, p_k) \in \Delta_k$ and $p_l = \frac{d}{k} \text{tr}(\rho_l \tau)$, where $\tau := (d+1) \sum_{j=1}^k p_j \rho_j - I$, $\tau \in \mathcal{P}(\mathbb{C}^d)$.

$$\Delta_k^{(2)} := \left\{ P \in \Delta_k \mid \sum_{j=1}^k p_j^2 = \frac{2d}{k(d+1)} \right\}.$$

This value is referred to as the **index of coincidence** and denoted $IC(P)$. Then [Harremoes&Topsoe '01]:

$$\arg \max_{P \in \Delta_k^{(2)}} H(P) = \left\{ s \left(p, \frac{1-p}{k-1}, \dots, \frac{1-p}{k-1} \right) \mid s \in S_k \right\},$$

$$\text{where } p = \frac{1}{k} \left(1 + \sqrt{\frac{(d-1)(k-1)}{d+1}} \right).$$

and:

$$\arg \min_{P \in \Delta_k^{(2)}} H(P) = \left\{ s(p, \underbrace{q, \dots, q}_m, \underbrace{0, \dots, 0}_{k-m-1}) \mid s \in S_k \right\},$$

$$\text{where } p = \frac{1-\sqrt{r_m}}{m+1}, q = \frac{m+\sqrt{r_m}}{m(m+1)}, m = \lfloor (IC(P))^{-1} \rfloor \text{ and } r_m = \frac{IC(P) - \frac{1}{m+1}}{\frac{1}{m} - \frac{1}{m+1}}.$$

For a 2-design POVM Π we have $\Delta_\Pi \subset \Delta_k^{(2)}$ and thus we get the following bounds on H_Π :

$$\min_{P \in \Delta_k^{(2)}} H(P) \leq \min_{P \in \Delta_\Pi} H(P) = \min_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_\Pi(\tau)$$

$$\max_{P \in \Delta_k^{(2)}} H(P) \geq \max_{P \in \Delta_\Pi} H(P) = \max_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_\Pi(\tau)$$

For which 2-design POVMs these bounds are met?

In other words:

for which 2 design POVM Π there exist $\tau_{\min}, \tau_{\max} \in \mathcal{P}(\mathbb{C}^d)$ such that

$$P(\tau_{\min}) \in \arg \min_{P \in \Delta_k^{(2)}} H(P)$$

and

$$P(\tau_{\max}) \in \arg \max_{P \in \Delta_k^{(2)}} H(P).$$

Theorem

Let Π be a 2-design POVM in \mathbb{C}^d . Then

$$\max_{P \in \Delta_k^{(2)}} H(P) = \max_{P \in \Delta_\Pi} H(P) = \max_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_\Pi(\tau)$$

if and only if Π is a **SIC-POVM**.

Sketch of the proof.

- For a SIC-POVM $\Pi = \{\frac{1}{d}\rho_j\}_{j=1}^{d^2}$ take $\tau = \rho_j$.
- To see the converse it is enough to show that the equality holds iff $k = d^2$.
- Use the fact that $\sum_{j=2}^k \rho_j = \frac{k}{d}I - \rho_1$ and the condition $\frac{k}{d}\rho_l = (d+1)\sum_{j=1}^k p_j \text{tr}(\rho_j \rho_l) - 1$ for $l = 1, \dots, k$.



Known examples for which the lower bound **is** achieved:

- the unique (tetrahedral) **SIC-POVM** in dimension **2**
 - minimal configuration: 'twin' SIC
- the **supersymmetric SIC-POVM** in dimension **3** (the **Hesse configuration**)
 - minimal configuration: complete set of MUB
- a **generic SIC-POVM** in dimension **3**
 - minimal configuration: an orthonormal basis
- the **complete set of MUB** in dimension **3**
 - minimal configuration: the supersymmetric SIC
- the **supersymmetric SIC-POVM** in dimension **8** (the **Hoggar lines**)
 - minimal configuration: 'twin' SIC

Known examples for which the lower bound is **not** achieved:

- SIC-POVM in dimension 4
- complete sets of MUBs in even dimensions
- configuration of 45 vectors in \mathbb{C}^5 (Example 18 in [Hoggar '82])

Note that for all known examples with attained lower bound we have:

$$IC(P)^{-1} = \frac{k(d+1)}{2d} \in \mathbb{N}$$

If $m := (IC(P))^{-1} = \frac{k(d+1)}{2d} \in \mathbb{N}$, then

$$\arg \min_{P \in \Delta_k^{(2)}} H(P) = \left\{ \underbrace{s\left(\frac{1}{m}, \dots, \frac{1}{m}\right)}_m, \underbrace{0, \dots, 0}_{k-m} \mid s \in S_k \right\}.$$

Let $\Pi = \{\frac{d}{k}\rho_j\}_{j=1}^k$ be a 2-design POVM in \mathbb{C}^d and let $\frac{k(d+1)}{2d} \in \mathbb{N}$. Then the following conditions are equivalent:

- $\min_{P \in \Delta_k^{(2)}} H(P) = \min_{P \in \Delta_\Pi} H(P) = \min_{\tau \in \mathcal{P}(\mathbb{C}^d)} H_\Pi(\tau)$
- There exists $J \subset \{1, 2, \dots, k\}$ such that $|J| = \frac{k(d-1)}{2d}$ and $\frac{2d}{k} \sum_{j \in J} \rho_j$ is an orthogonal projection onto $(d-1)$ -dimensional subspace of \mathbb{C}^d .
- Let $\rho_j = |\psi_j\rangle\langle\psi_j|$ for some $\psi_j \in \mathbb{C}^d$. Then there exists $J \subset \{1, 2, \dots, k\}$ such that $|J| = \frac{k(d-1)}{2d}$, $\dim(\text{span}\{\psi_j\}_{j \in J}) = d-1$ and $\{\psi_j\}_{j \in J}$ is a **normalized tight frame** in $\text{span}\{\psi_j\}_{j \in J}$.
- There exists $J \subset \{1, 2, \dots, k\}$ such that $|J| = \frac{k(d-1)}{2d}$ and $\sum_{j \in J} p_j(\tau) \leq \frac{1}{2}$ for all $\tau \in \mathcal{P}(\mathbb{C}^d)$.

Quantum bayesianism [Appleby, Fuchs, Stacey & Zhu, 2016]

- "Bureau of Standards" measurement: SIC-POVM $E = \{E_j\}_{j=1}^{d^2}$
- an arbitrary POVM $F = \{F_j\}_{j=1}^m$
- let $r(j|i) = \text{tr}(E_i F_j)/d$
- for $\tau \in \mathcal{S}(\mathbb{C}^d)$ let $p_\tau(j) = \text{tr}(\tau E_j)$ and $q_\tau(j) = \text{tr}(\tau F_j)$.
- **URGLEICHUNG:**

$$q_\tau(j) = \sum_{i=1}^{d^2} (\alpha p_\tau(i) - \beta) r(j|i)$$

with $\alpha = d + 1$ and $\beta = 1/d$.

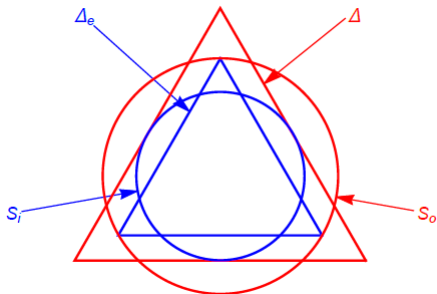
- **QPLEX** = $\{(p_\tau(1), \dots, p_\tau(d^2)) | \tau \in \mathcal{S}(\mathbb{C}^d)\}$
- BUT for any 2-design POVM $E = \{E_j\}_{j=1}^k$ and $r(j|i) = \text{tr}(E_i F_j)(d/k)$ we also have

$$q_\tau(j) = \sum_{i=1}^k (\alpha p_\tau(i) - \beta) r(j|i)$$





with $\alpha = d + 1$ and $\beta = d/k$

Properties of qplex:

- contained in outer ball, contains inner ball
- the same as quantum state space
- Can geometrical properties and abstract definitions be transferred into general 2-design POVM case?



- How general qplex can be located in the large simplex?
- Can this better understanding of the structure of Δ_{Π} lead to progress in the minimal entropy problem?

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