

# ON Maximal Entanglement

CODES, HOLOGRAPHY, QMA

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# Quantum Information is bringing new insights

Traditional emphasis on operators



$$H |\psi\rangle = E |\psi\rangle$$



QI emphasis on states

## Example: Quantum Phase Transitions

$H$

Criticality  
RG flows on coupling constants  
OPE, Conformal Symmetry

$|\psi\rangle$

Scaling of entropy  
RG flows on states  
Distribution of entanglement: MERA

Von Neumann entropy

$$\rho_A = \text{Tr}_B |\psi\rangle_{AB} \langle \psi|$$

$$S(\rho_A) = -\text{Tr} \rho_A \log \rho_A$$

*Fixed points: scaling with the size  $L$*

$$S(\rho_A) = \frac{c}{3} \log L \ll L$$

Callan-Wilczek 94, Vidal-jil-Rico-Kitaev 02

$$S \sim n - \text{const}$$

Random states, QMA problems,  
Local translational invariant higher d

$$S \sim .8858 n + \text{const}$$

Prime state

$$S \sim n^{\frac{d-1}{d}} + \text{const}$$

Area law in d-dimensions

$$S \sim \frac{c}{3} \log n + \text{const}$$

Critical scaling in d=1

$$S \sim \log(\xi) = \text{const}$$

Finitely correlated states  
away from criticality

# Outline

Introduction

Absolute Maximal Entanglement

H-code

The most entangled spin chain

Gauge theory from MaxEnt

# Absolute Maximal Entanglement

Def: Absolute Maximal Entanglement

$$|\psi\rangle \in H_d^{\otimes n}$$

$$\rho_m = \text{Tr}_{n-m} |\psi\rangle\langle\psi| = \frac{1}{d^m} I_{d^m} \quad \forall m$$

# Absolute Maximum Entanglement

Helwig, Cui, JIL, Riera, Lo (2012)

Multipartite teleportation

Quantum secret sharing

Reed-Solomon codes

Helwig thesis (2014)

## Holography

JIL, Sierra (2015); Pastawski, Yoshida, Harlow, Preskill (2015)

## Combinatorial design

Goyneche, Alsina, JIL, Riera, Zyczkowsky (2015)

Previous work on perfect states, balanced states, LU, LOCC,...



## 4-qubit maximally entangled state?

**NO: there is no 4-qubit state maximally entangled in all its partitions**

Monogamy of entanglement vs local dimensionality

Scott, Facchi, Parisi, ...

AME exist for  $n=2,3,5,6$  qubits

AME do not exist for  $n=4$  and  $n>7$

AME for  $n=7$  is not known, though numerical evidence is against it

# 4-qutrit Absolute Maximally Entangled state in all partitions

- First non-trivial generalization of EPR states
- Symmetric (mod relabeling)
- Higher local dimensionality leaves space for entanglement
- Maximal surprise = Error correction

$$|P\rangle_{ABCS} = \frac{1}{3} \sum_{i,j=0,1,2} |i, j, i+j, i+2j\rangle$$

$$\rho_{AB} = \rho_{AC} = \rho_{AS} = \frac{1}{9} I_9$$

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$|1\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$|2\rangle = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

mod 3 notation

$$|AME(4,3)\rangle_{ABCS} = \sum_{i,j=0,1,2} |i, j, i+j, i+2j\rangle$$

$i$	$j$	$i+j$	$i+2j$
0	0	0	0
0	1	1	2
0	2	2	1
1	0	1	1
1	1	2	0
1	2	0	2
2	0	2	2
2	1	0	1
2	2	1	0



**basis** → **basis**

same for any partition

$$t_{abcs} = A_{abc}^s$$



$s, a$  fixed by  $b, c$

Hamming distance = 3

(error correction code, Reed-Solomon 1950)

# AME and Reed-Solomon codes

$$c_i = e_i \cdot G$$

↙ ↘

codewords      basis

$$AME(4,3) \quad G = \begin{pmatrix} 1 & 1 & 1 & 0 \\ 0 & 1 & 2 & 1 \end{pmatrix}$$

$$AME(n,d) \quad G = \begin{pmatrix} 1 & 1 & \dots & 1 & 0 \\ a_1 & a_2 & \dots & a_q & 0 \\ \cdot & \cdot & \dots & \cdot & 0 \\ a_1^k & a_2^k & \dots & a_q^k & 1 \end{pmatrix}$$

$q = d = \text{Prime}$

$n = d + 1 \quad k = n/2$

# Multi-Unitarity

$i$	$j$	$i+j$	$i+2j$
0	0	0	0
0	1	1	2
0	2	2	1
1	0	1	1
1	1	2	0
1	2	0	2
2	0	2	2
2	1	0	1
2	2	1	0

$$t_{abcd} = U_{(ab)(cd)} = U_{(ac)(bd)} = U_{(ad)(bc)}$$

$$|AME\rangle = \sum_{i_1, \dots, i_m} |i_1, \dots, i_m\rangle U_{(i_1, \dots, i_m), (i_{m+1}, \dots, i_n)} |i_{m+1}, \dots, i_n\rangle$$

$$U U^\dagger \sim I, \text{ isometry}$$

## AME support

Local Unitary (LU) invariance allows for a minimal support

$AME(4,3)$  needs only 9 superposed elements

$AME(6,2)$  needs a minimum of 16 superposed elements

# AME

Are AME related to any maximal violation of Bell Inequalities?

Are AME related to some underlying entanglement structure of physical systems?

Is Maximal Entropy possible in Condensed Matter systems?

# Holographic codes



# Tensor Networks

Simulation of quantum systems  
on a classical computer

Exact Simulation

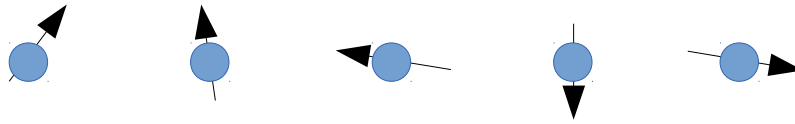
Monte Carlo

**Tensor Networks**

## Tensor Networks

- Variational approach
- Escape Sign problem
- Support area law entanglement
- Time-evolution
- Libraries: DMRG, iTEBD,..
- Include symmetries

# Example: Matrix Product State



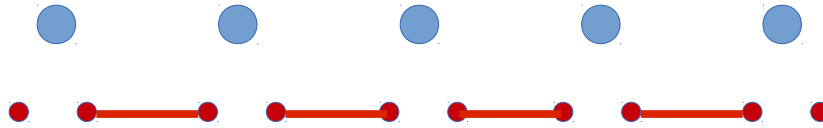
$$|\psi\rangle = \sum_{i_1 \dots i_n} c^{i_1 \dots i_n} |i_1 \dots i_n\rangle \in (\mathbb{C}^2)^{\otimes n}$$

$2^n$  coefficients



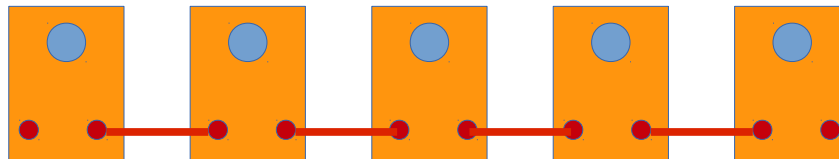
$$c^{i_1 \dots i_n} = \text{Tr} A^{i_1} \dots A^{i_n}$$

Populate the system with entangled pairs



$$|\varphi\rangle = \sum_{\alpha=1, \dots, \chi} |\alpha \alpha\rangle$$

and project ancillary states onto physical indexes



Just capture correlations!!

$$A_{\alpha\beta}^i$$

$$c^{i_1 \dots i_n} = \text{Tr} A^{i_1} \dots A^{i_n}$$

Exponential Kolmogorov complexity reduction!

$d^n$  coefficients  $\longrightarrow$   $nd\chi^2$  coefficients

Adaptive algorithm to entanglement

$$\chi = 1$$

Product state

$$\chi = \text{poly}(n)$$

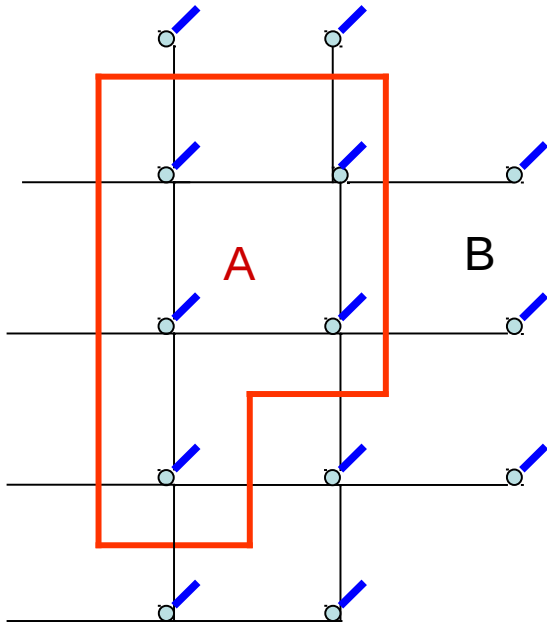
Slight entanglement

1D quantum phase transitions

$$\chi = 2^{\frac{n}{2}}$$

Maximal entanglement

Problems in QMA

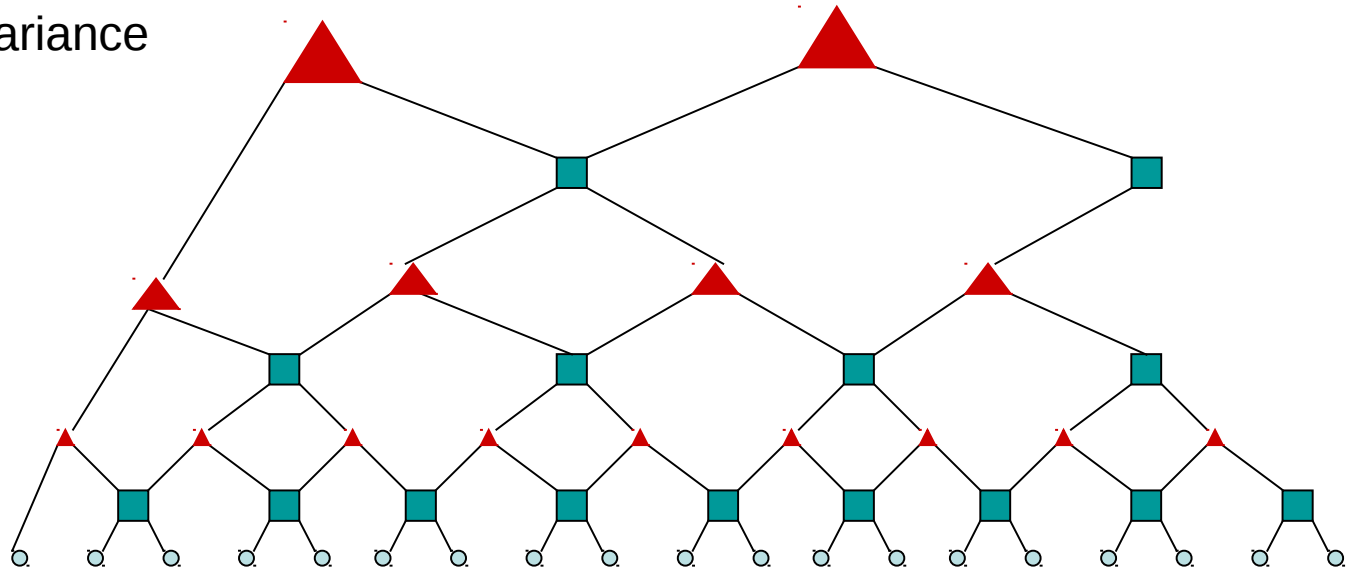


Support for entanglement entropy  
is proportional to the boundary

$$S_A = S_B \propto L$$

**Area law**

- |      |       |                          |
|------|-------|--------------------------|
| MPS  | D=1   | translational invariance |
| PEPS | D=2,3 | translational invariance |
| Tree |       | scale invariance         |
| MERA |       | scale invariance         |

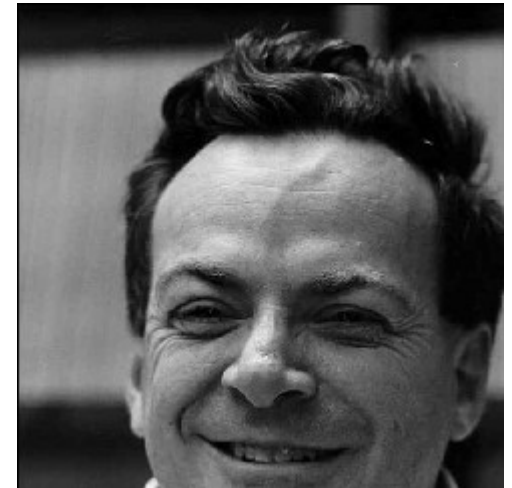


Cirac, Verstraete  
Vidal

Tensor Networks are compressors of correlations

$$|\psi\rangle = \sum_{i_1 \dots i_4 = 1, \dots, 9} c^{i_1 \dots i_4} |i_1 \dots i_4\rangle \in (C^9)^{\otimes 4}$$

↑                      ↑  
level of grey      pixel address



$$c^{i_1 \dots i_n} = \text{Tr } A^{i_1} \dots A^{i_n}$$

$$\dim A = \chi$$

Max  $\chi = 81$



$\chi=1$   
error=27

$\chi=2$   
error=14.8

$\chi=3$   
error=9.2

$\chi=4$   
error=7.1

```
graph TD; A([Tensor Networks]) --> B([simulation of Hamiltonians]); A --> C([entanglement structure of States]);
```

Tensor Networks

simulation  
of Hamiltonians

entanglement structure  
of States

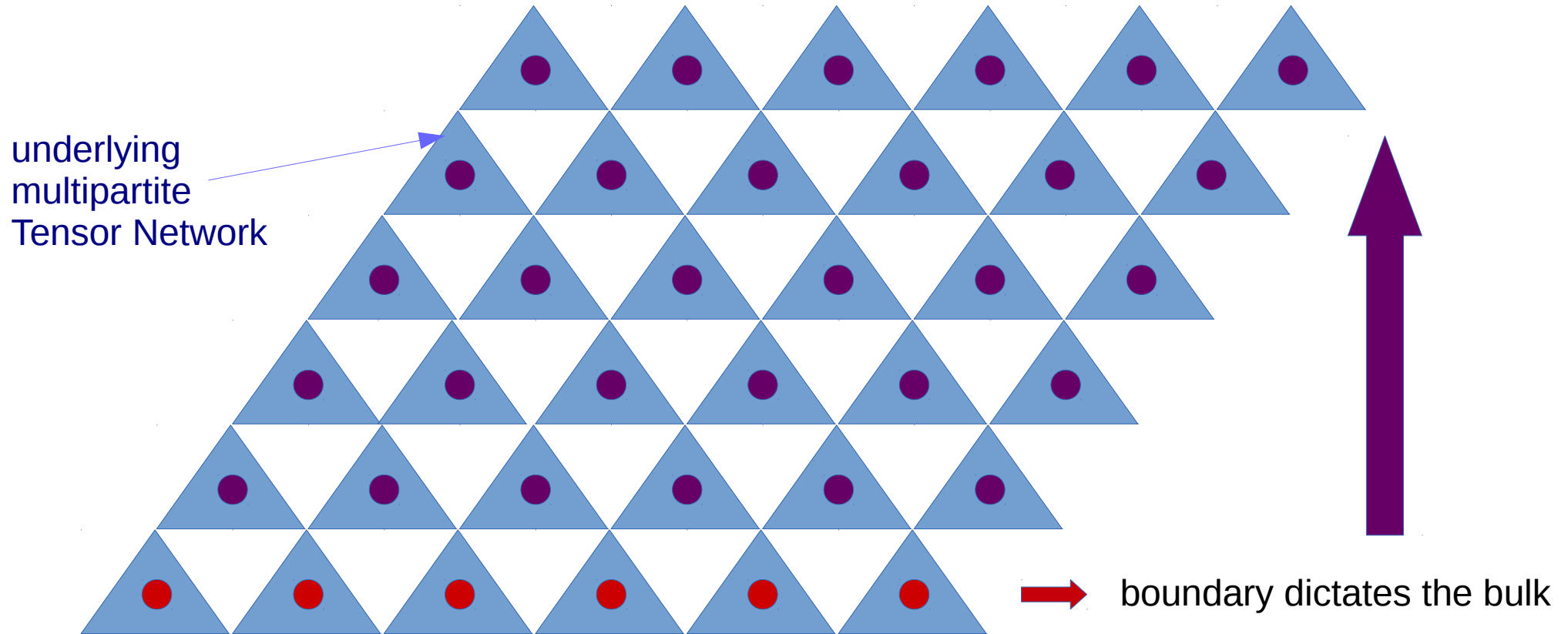
## H code

$$|H\rangle = \sum_i |i\rangle_{boundary} |\varphi_i\rangle_{bulk}$$

$$\langle i|j\rangle = \langle \varphi_i|\varphi_j\rangle = \delta_{ij}$$

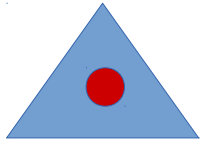
$|i\rangle, |\varphi_i\rangle$  product states

# H code



$$|H\rangle = \sum_i |i\rangle_{\text{boundary}} |\varphi_i\rangle_{\text{bulk}}$$

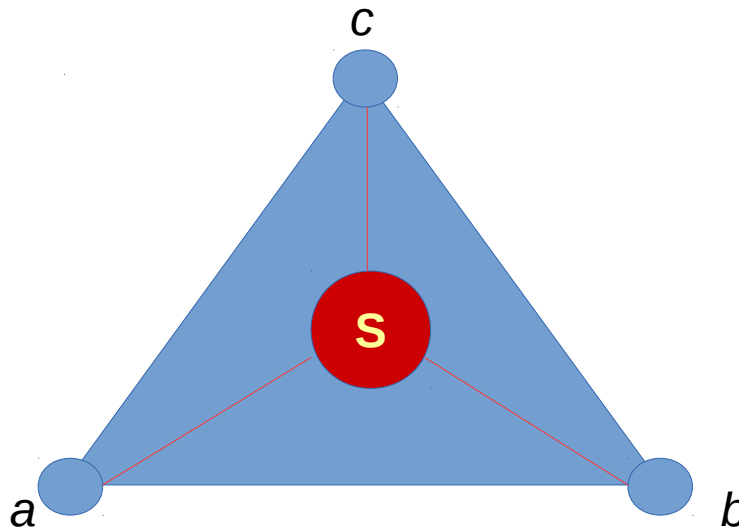




Choice of tensor network = choice of simplex + connection of simplices

Choice of simplex

$$|\psi\rangle_{ABCS} = \sum_{abcs} t_{abcs} |abc\rangle_{\text{ancillae}} |s\rangle_{\text{physical}}$$

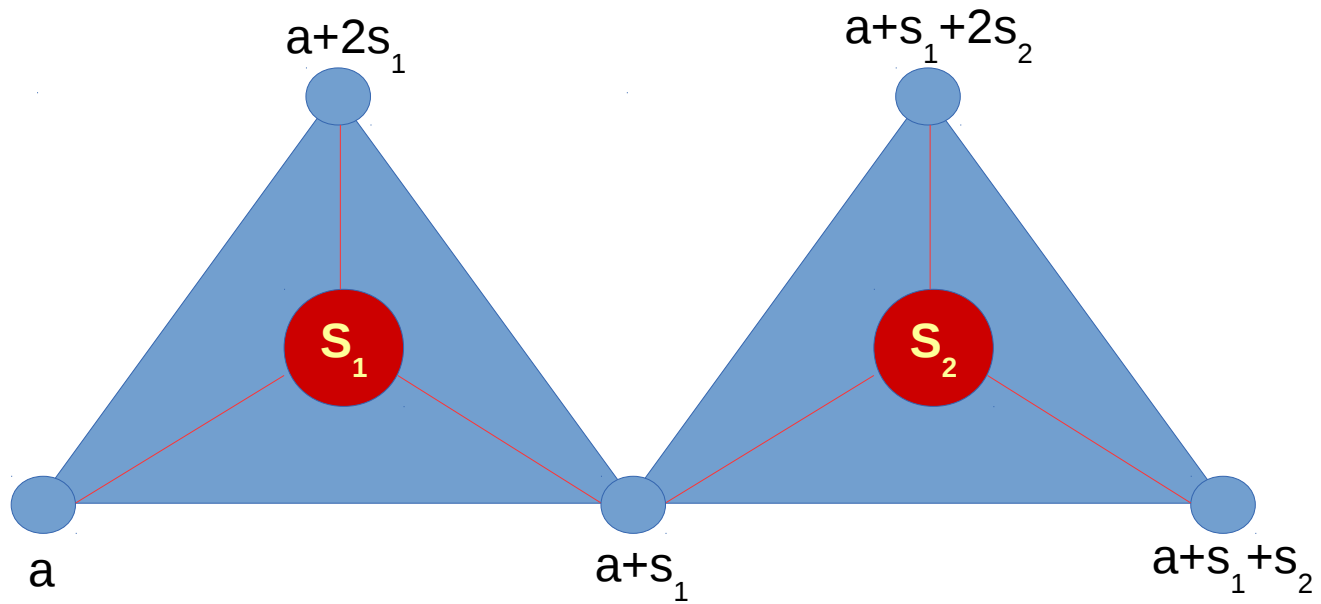


$$t_{abcs} = A_{abc}^s$$

Let's combine

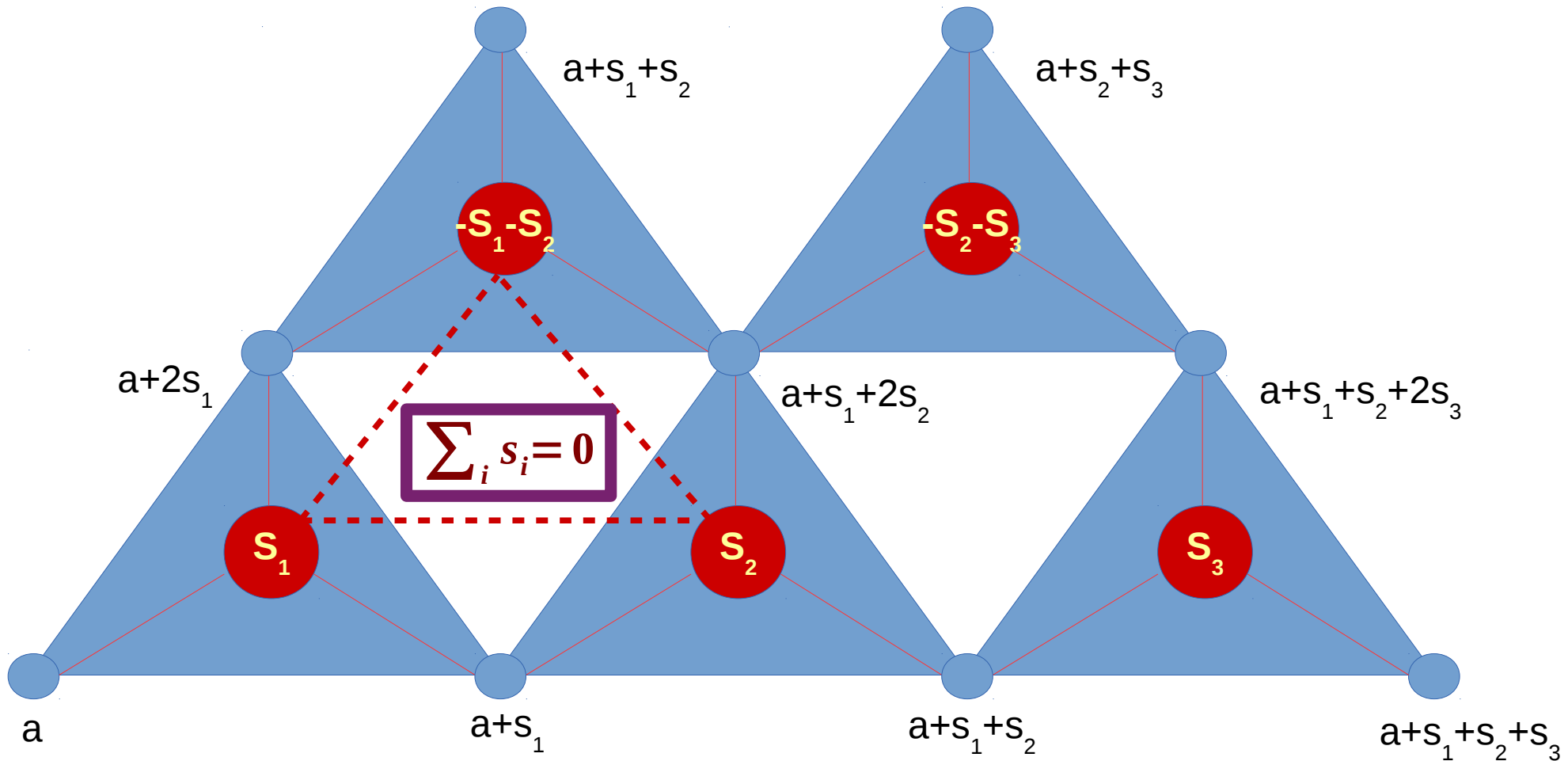
Tensor Network + AME

## Build network by linking ancillae



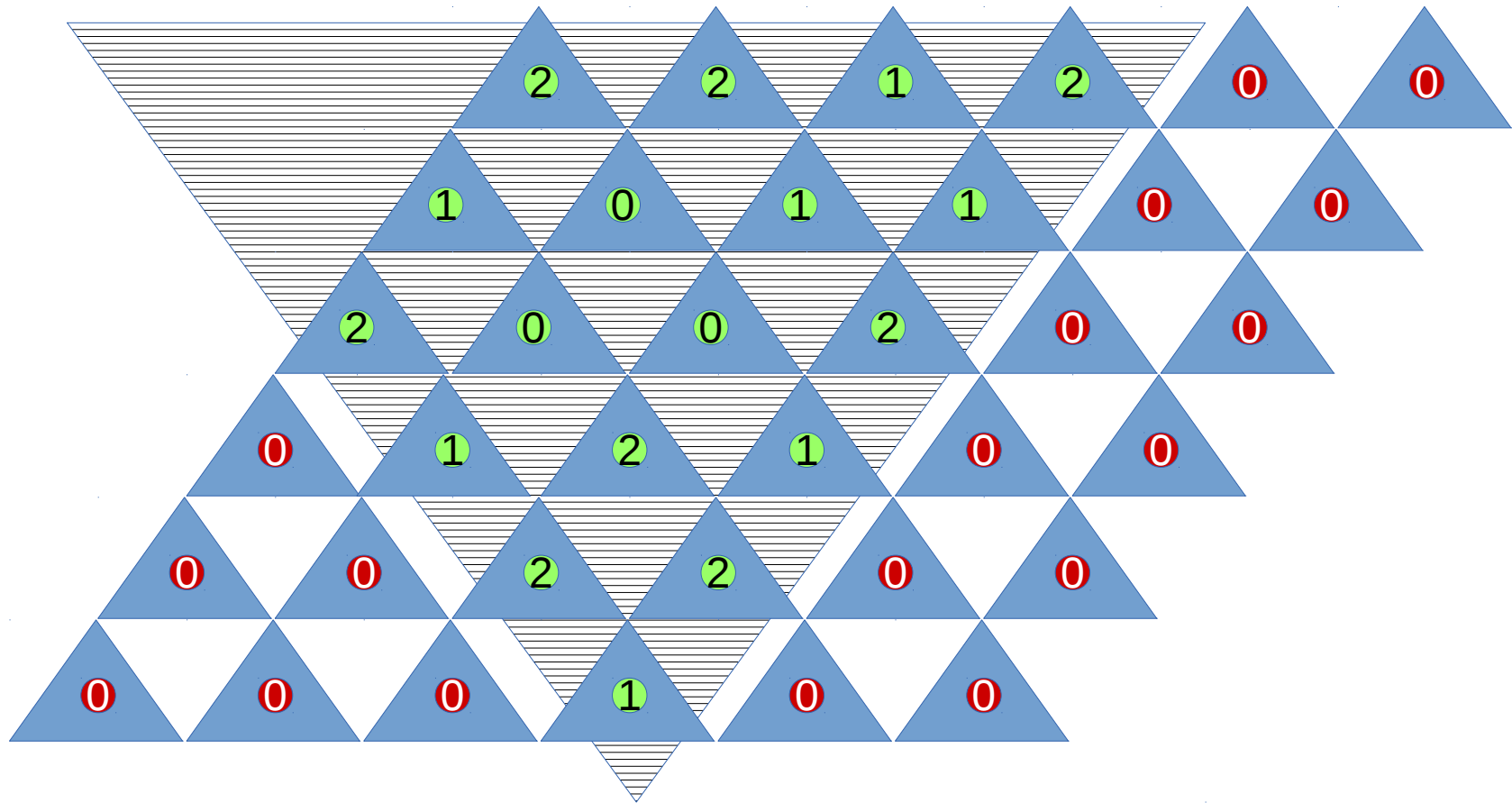
Ancillae are shared from one simplex to the contiguous ones

Emergence of **neutralization rule** = holography



Neutralization = Gauss law ?

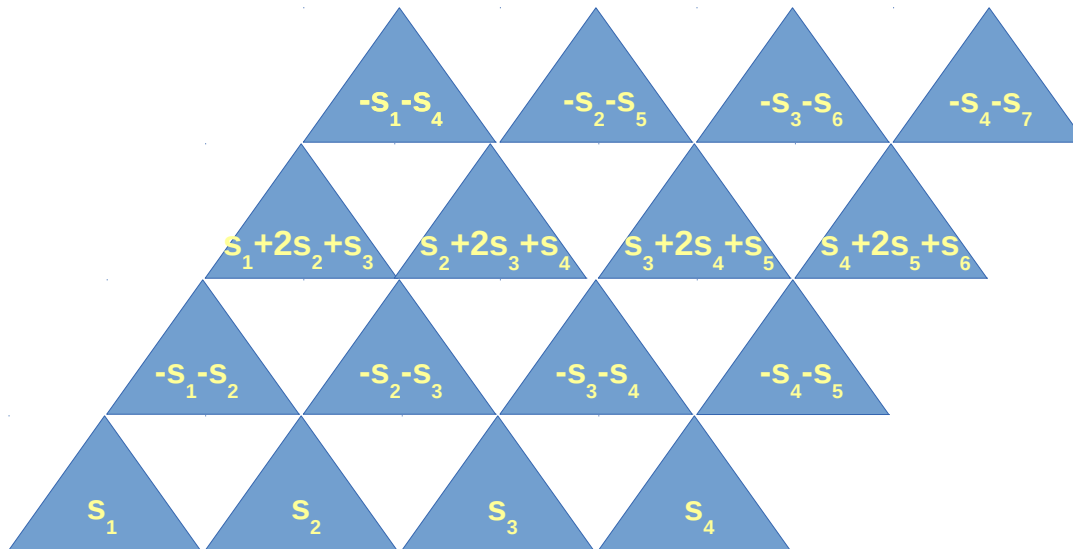
# Light-cone influence of a boundary qutrit



# States on a Torus

The neutralization rule may enter in conflict with periodicity

Yet, if  $s_1 = s_4$  then  $-s_1 - s_4 = -2s_1 = s_1$



Periodicity shows up for size  
 $n=m=3^k$

and for special sizes  
 $(n,m)=(5,40),(7,182),(11,121),\dots$

A proof of periodicity for  $(n,m)=3^k$

$$\begin{pmatrix} s_1' \\ s_2' \\ \dots \\ s_n' \end{pmatrix} = \begin{pmatrix} -1 & -1 & 0 & \dots & 0 \\ 0 & -1 & -1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ -1 & 0 & 0 & \dots & -1 \end{pmatrix} \begin{pmatrix} s_1 \\ s_2 \\ \dots \\ s_n \end{pmatrix}$$

$$T_n = -I - U_n$$

$$U_n = \begin{pmatrix} 0 & 1 & 0 & \dots & 0 \\ 0 & 0 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & 0 & 0 & \dots & 0 \end{pmatrix}$$

$$U_n^n = I$$

$$T_n^n = (-I - U_n)^n \pmod{3} = - \sum_{r=0}^n \binom{n}{r} U_n^r \pmod{3} = -(I + U_n^n) \pmod{3} = -2I \pmod{3} = I$$

$n = 3^k$

(Some tori sizes only accept a subset of states)

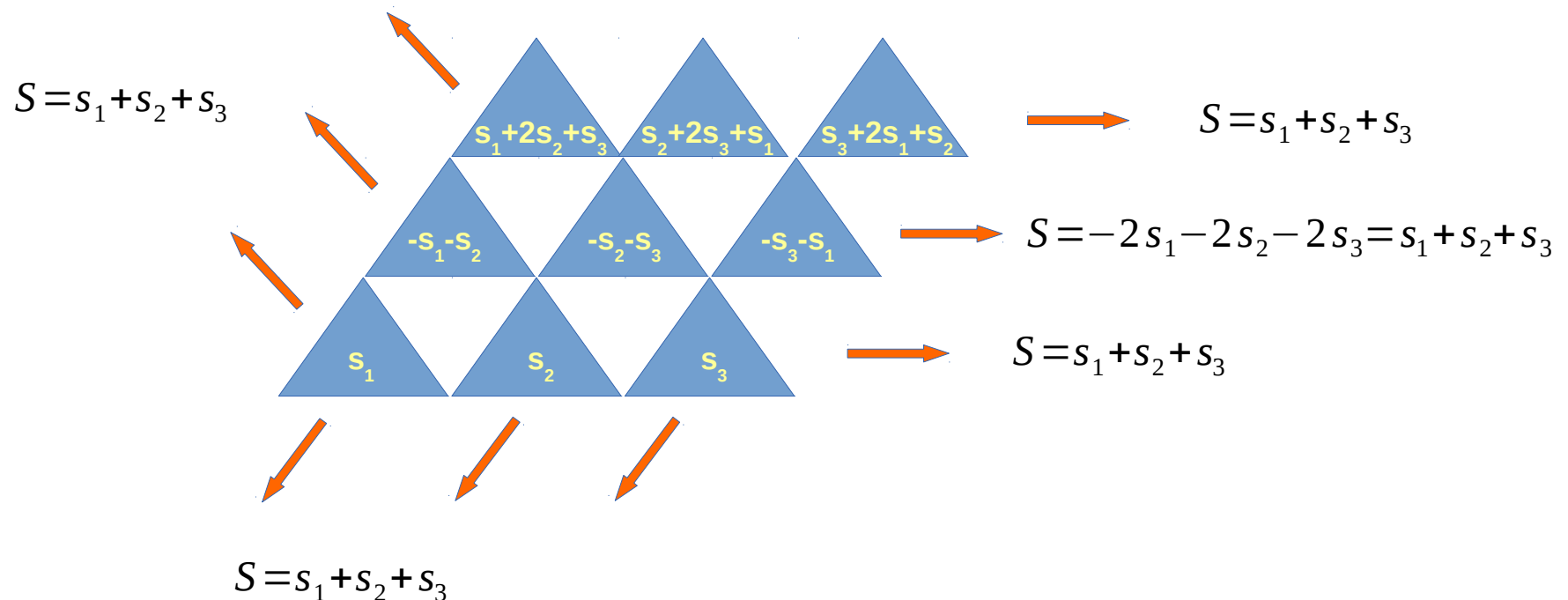
# Topological properties of the 3x3 torus

$$S = \sum_i s_i$$

$$Q = \prod_i z_i$$

$$z_i = q^{1+s_i} \quad q = e^{\frac{2\pi i}{3}}$$

REDUNDANCY: **All cycles carry the same redundant information !**



Same is valid for larger tori



## H code

$$|H_S\rangle = \sum_{s_i / \sum s_i = S} |s_i\rangle_{boundary} |f(s_i)\rangle_{bulk}$$

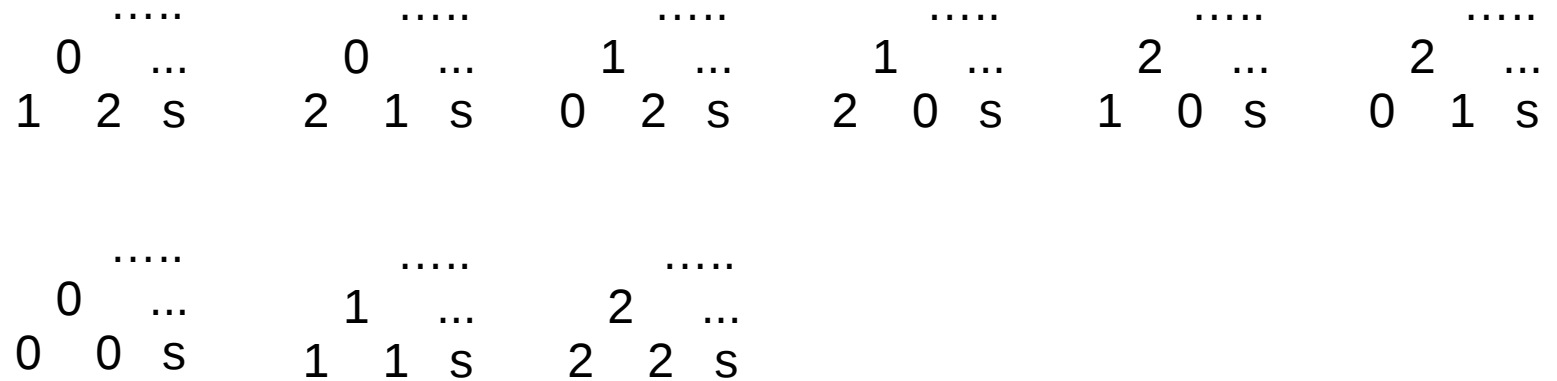
A qutrit can be encoded in  $(|H_0\rangle, |H_1\rangle, |H_2\rangle)$

through the value of topological redundant observables

If Hamming distance is large, robust quantum memory is possible

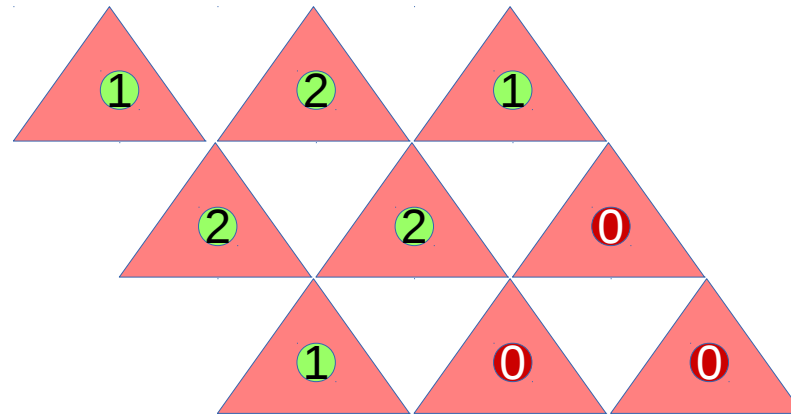
## Symmetry + Neutralization rule

Any triangle will carry all possibilities equally weighted



**No local observable can distinguish the three phases**

Revisit the light-cone in 3x3 torus: large Hamming distance!



Hamming distance = 6

For size  $n=m=3^k$  Hamming distance =  $6^k$

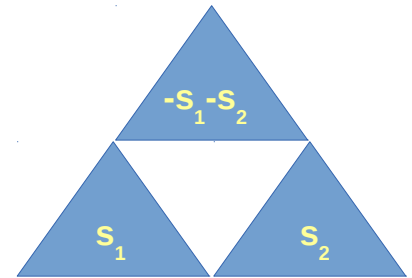
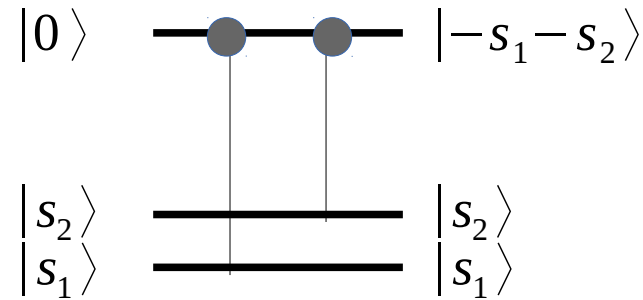
$dist(3 \times 3) = 6$      $dist(9 \times 9) = 36$

Fractal dimension

boundary =  $n$     volum =  $n^2$

$dist = n^{(\log 6 / \log 3)} = n^{1.63091}$

# Construction of a H-code as a cellular automata



bulk  
sequential  
neutralization gates

$$U_{neutralization} = U_{1,3} U_{2,3}$$

$$U|a\rangle|b\rangle = |a\rangle|ba^2\rangle$$

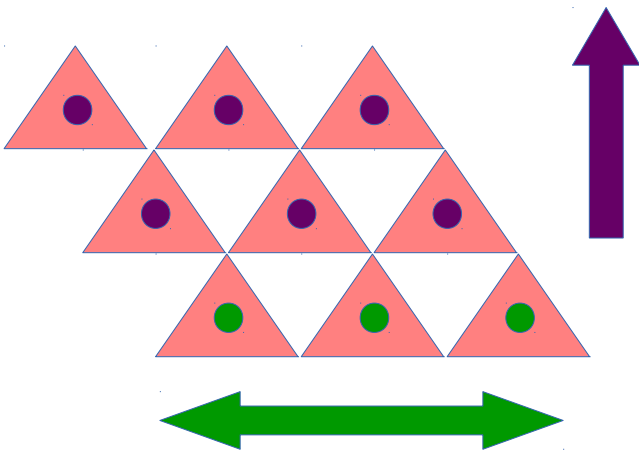
$$a, b = 1, q, q^2$$

$$q = e^{\frac{2\pi i}{3}}$$

$$H_{boundary} = - \sum_{i=1}^n \sum_{a+b=0} X_i^a X_{i+1}^b$$

$$X = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

protects the total sum of the row



boundary  
symmetrization

# Parent Hamiltonian

Trivial proof of non-locality of a parent Hamiltonian

$$\mathbf{H} |H\rangle = \mathbf{H} (|\psi_1\rangle + |\psi_2\rangle + \dots)$$

dist =  $6^k$

$$\mathbf{H} |\psi_1\rangle \sim |\psi_2\rangle \quad \longrightarrow \quad \mathbf{H} \text{ must be } 6^k\text{-local}$$

$$\mathbf{H} = \mathbf{H}_{Z, \text{neutralization}} + \mathbf{H}_{X, \text{link triangles}} \quad [\mathbf{H}_X, \mathbf{H}_Z] = 0$$

$$\mathbf{H}_{Z, \text{neutralization}} = \sum_{i, j, k \in \Delta} \left( 2 - Z_i Z_j Z_k - (Z_i Z_j Z_k)^2 \right) \quad Z = \begin{pmatrix} 1 & 0 & 0 \\ 0 & q & 0 \\ 0 & 0 & q^2 \end{pmatrix}$$

Produces a degenerate ground state, where every triangle is neutralized

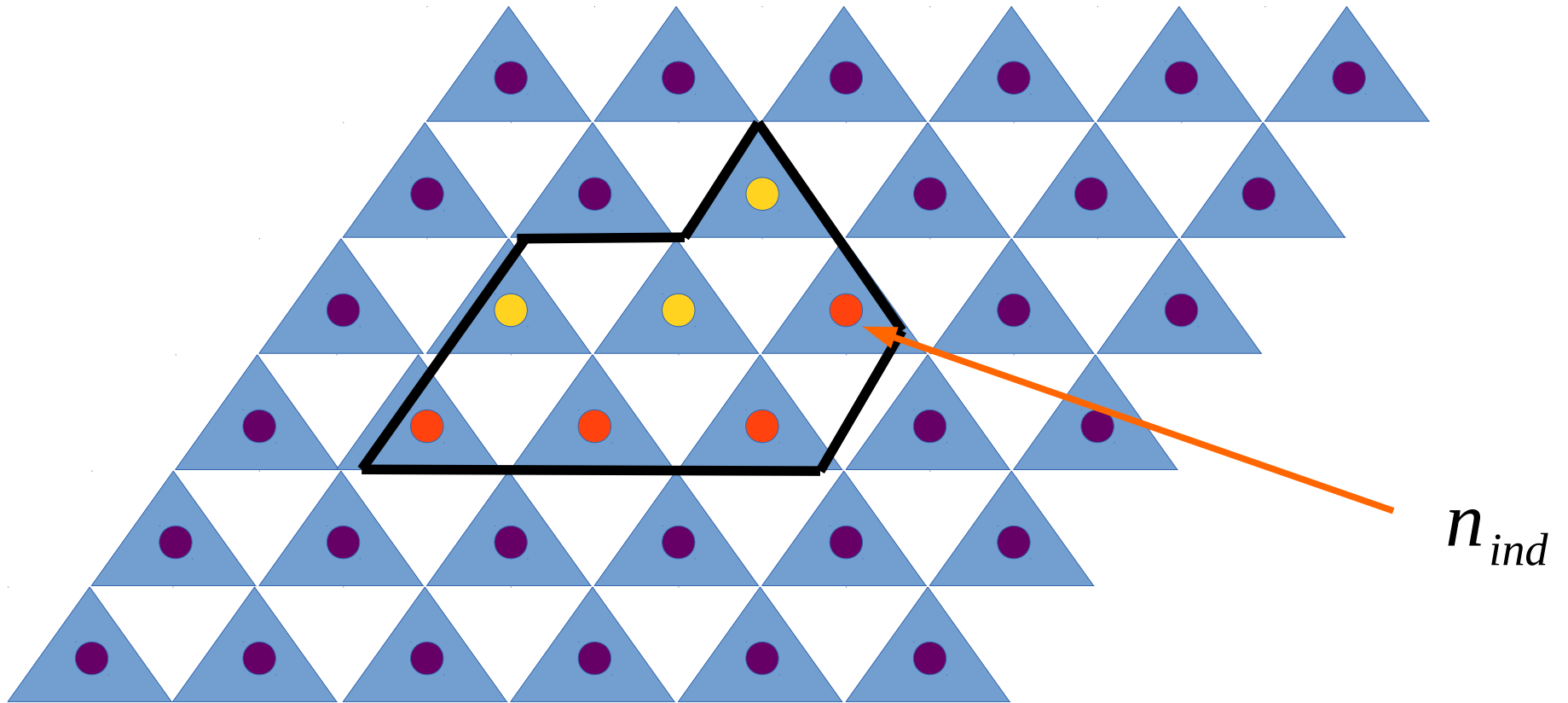
To break the degeneracy and obtain a single ground state, link triangles with

$$\mathbf{H}_{X, \text{link triangles}} = - \sum_{p, n=0,1,2} \prod_{a, b=1, \dots, 3^k} X_{a, b}^{-n(a-b+p)}$$

where  $X$  acts on the site  $a, b$

*$\mathbf{H}$  is not unique! The tensor network is not injective*

# Entropy of entanglement



$$S = -\text{Tr} \rho_i \log_3 \rho_i = n_{ind}$$

$$S_{top} = S_{ABC} - S_{AB} - S_{AC} - S_{BC} + S_A + S_B + S_C = -1$$

# Multi-Unitarity

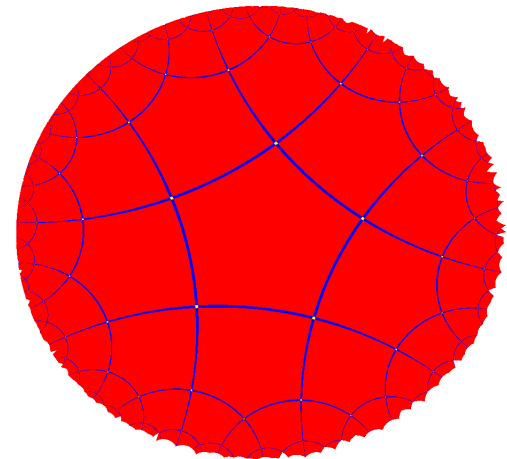
AME iff

$$t_{abcd} = U_{(ab)(cd)} = U_{(ac)(bd)} = U_{(ad)(bc)}$$

No sense of direction for RG interpretation  
Emergence of symmetry and holography

Is it necessary? K-uniformity?

Isometries: Pastawski, Yoshida, Harlow, Preskill: AME(5,2)  
Combinatorial design: Goyeneche, Alsina, JIL, Życzkowski





The most entangled spin chain

CFT

Entropy is maximal at criticality

$$S \sim \frac{c}{3} \log n + \text{const}$$

QMA

2-local translational invariant is QMA

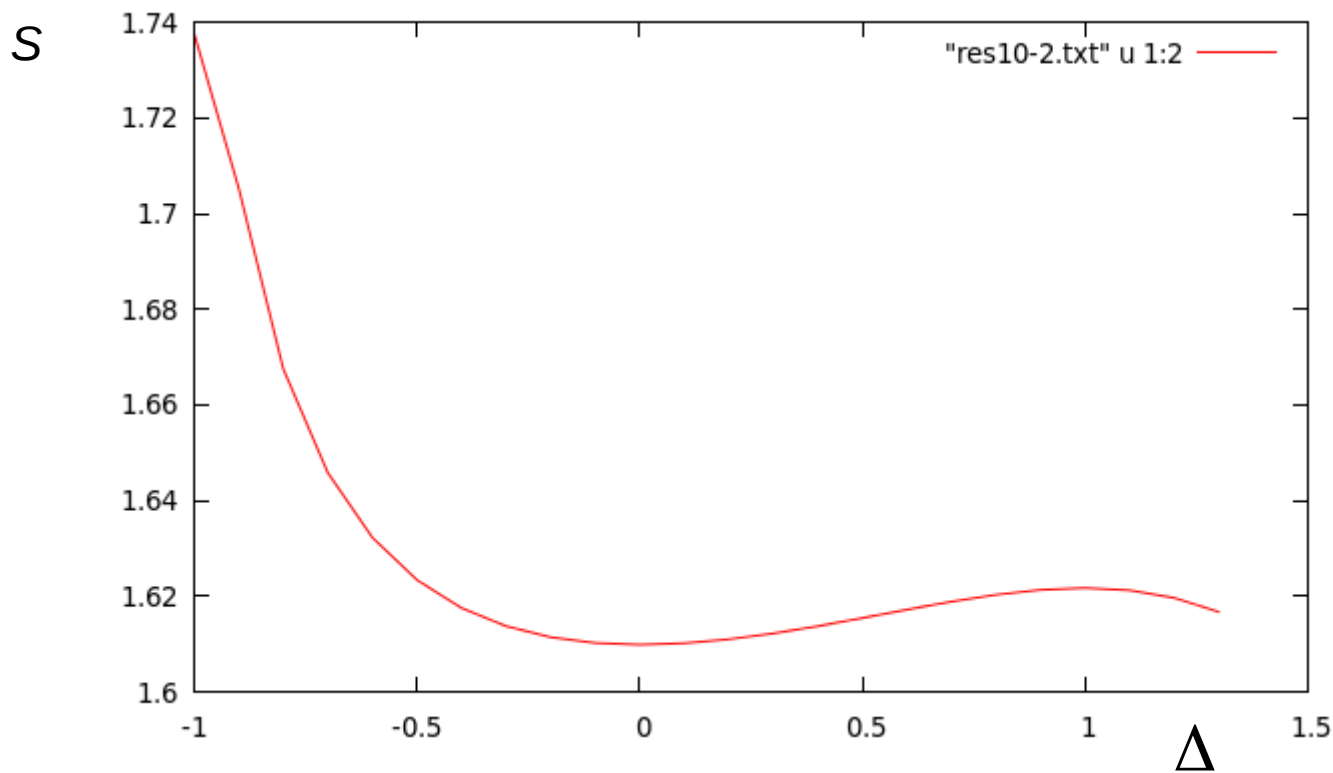
A spin 2 model proposed by Movassagh-Shor

$$S \sim \sqrt{n}$$

Minimal spin where QMA takes over?

Spin  $\frac{1}{2}$ : numerical evidence sets max entanglement along the line

$$H = \sum_i \left( \sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y + \Delta \sigma_i^z \sigma_{i+1}^z \right)$$



Ferromagnetic Heisenberg:

$$\Delta \rightarrow -1^+$$

$$S = \frac{1}{2} \log(n+1)$$

↖  
No CFT !

Frustration beats conformal symmetry!

Preliminary results for higher spins

$$H = \sum_{a \notin \text{Casimir}} \lambda_i^a \lambda_{i+1}^a + \Delta \sum_{a \in \text{Casimir}} \lambda_i^b \lambda_{i+1}^b$$

Maximum at  $\Delta \rightarrow -1^+$

Spin 1/2	Log n
Spin 1	Log n <sup>2</sup>
Spin 3/2	Log n <sup>3</sup>

No sign of QMA yet

# Gauge theory from MaxEnt

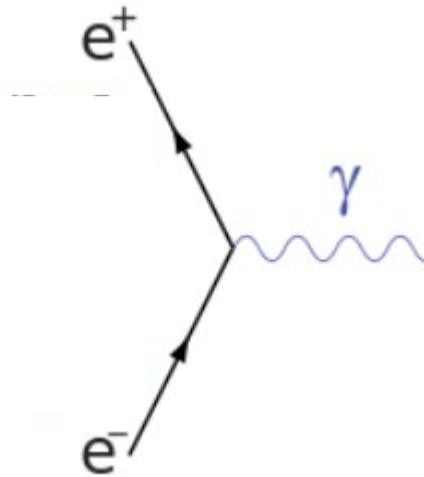
Conformal symmetry is obtained as a (local) max of entropy in 1D

Could we think of gauge symmetry as a consequence of maximal entanglement

**MaxEnt**

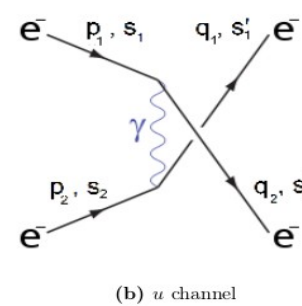
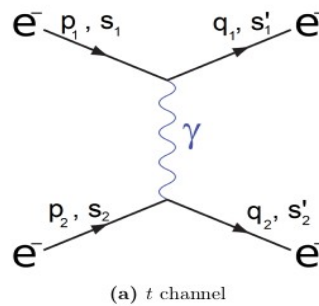
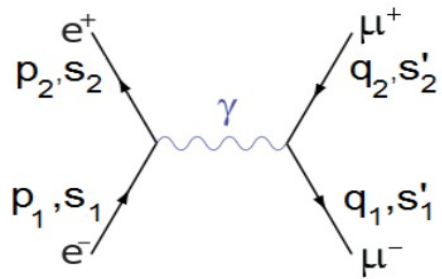
?

Let the QED vertex free



$$e \gamma^\mu \rightarrow e G^\mu$$

And fix  $G^\mu$  demanding max entanglement in processes



.....

- 1) Impose hermiticity, parity, charge conjugation and time reversal
- 2) Impose generation of MaxEnt and conservation of MaxEnt for Moller, Bhabha, Mott scatterings

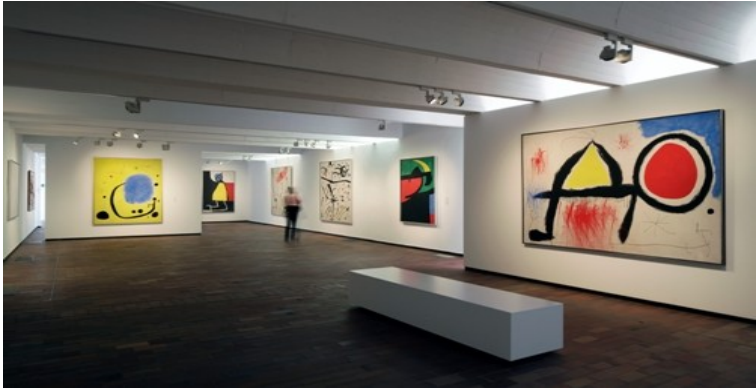
$$G^0 = \begin{pmatrix} 0 & \pm I \\ \mp I & 0 \end{pmatrix} \quad G^i = \begin{pmatrix} 0 & \pm \sigma^i \\ \mp \sigma^i & 0 \end{pmatrix}$$

Further signs may be constrained by Compton scattering



# CONCLUSION

There is life in the MaxEnt part of the Hilbert space



THANKS

see you in BCN



TNS ancillary maximal entanglement generates interesting states

Absolute Maximally Entangled 4-qutrit state generates a holographic code

3-fold degenerate ground state, large hamming distance, non-local observables

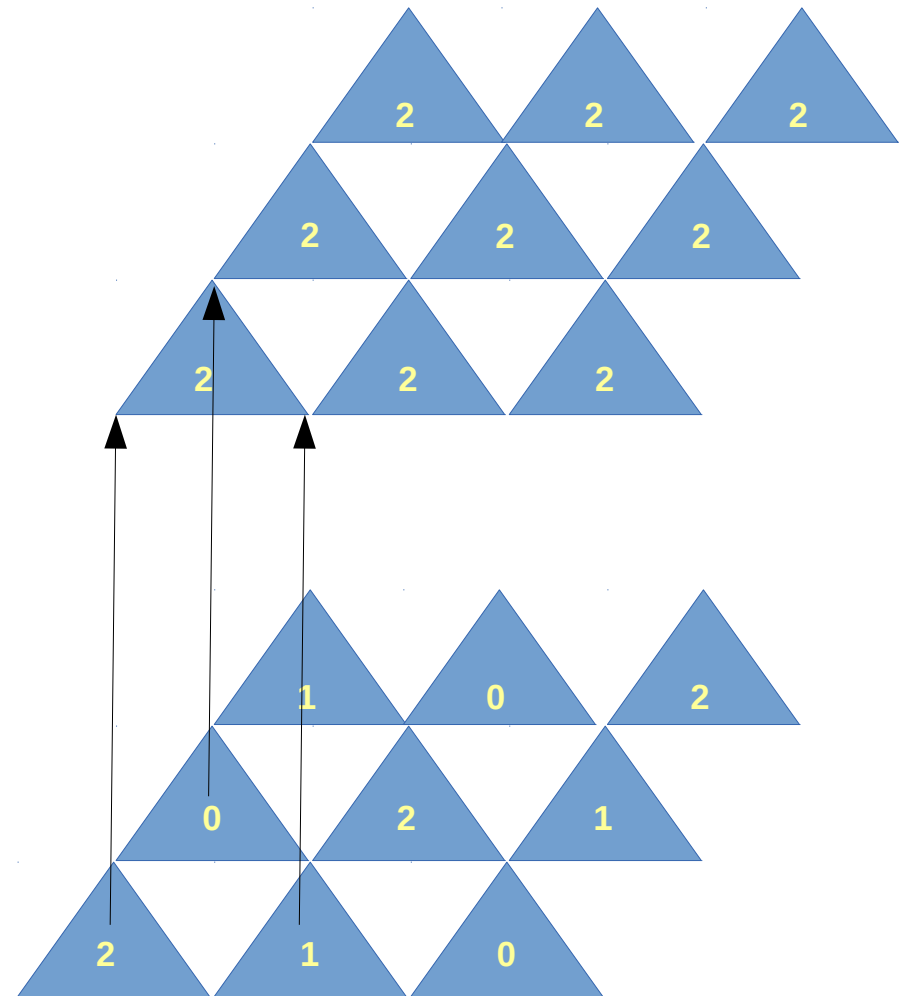
Holography = Multi-Unitarity ?

Gauss law is coming from gauge symmetry  
(Celi, Tagliacozzo, Sierra, JIL)

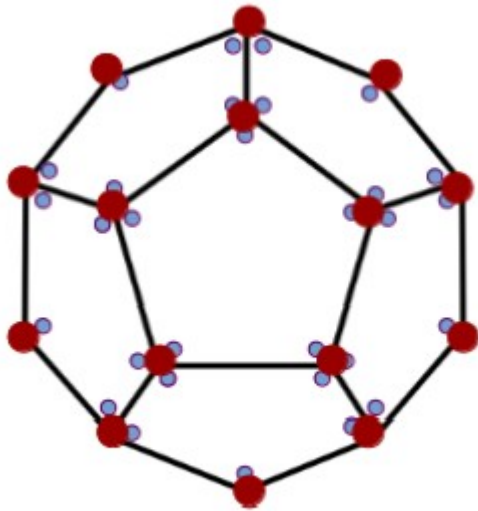
3D ?

2 non-trivial layers + fixed point

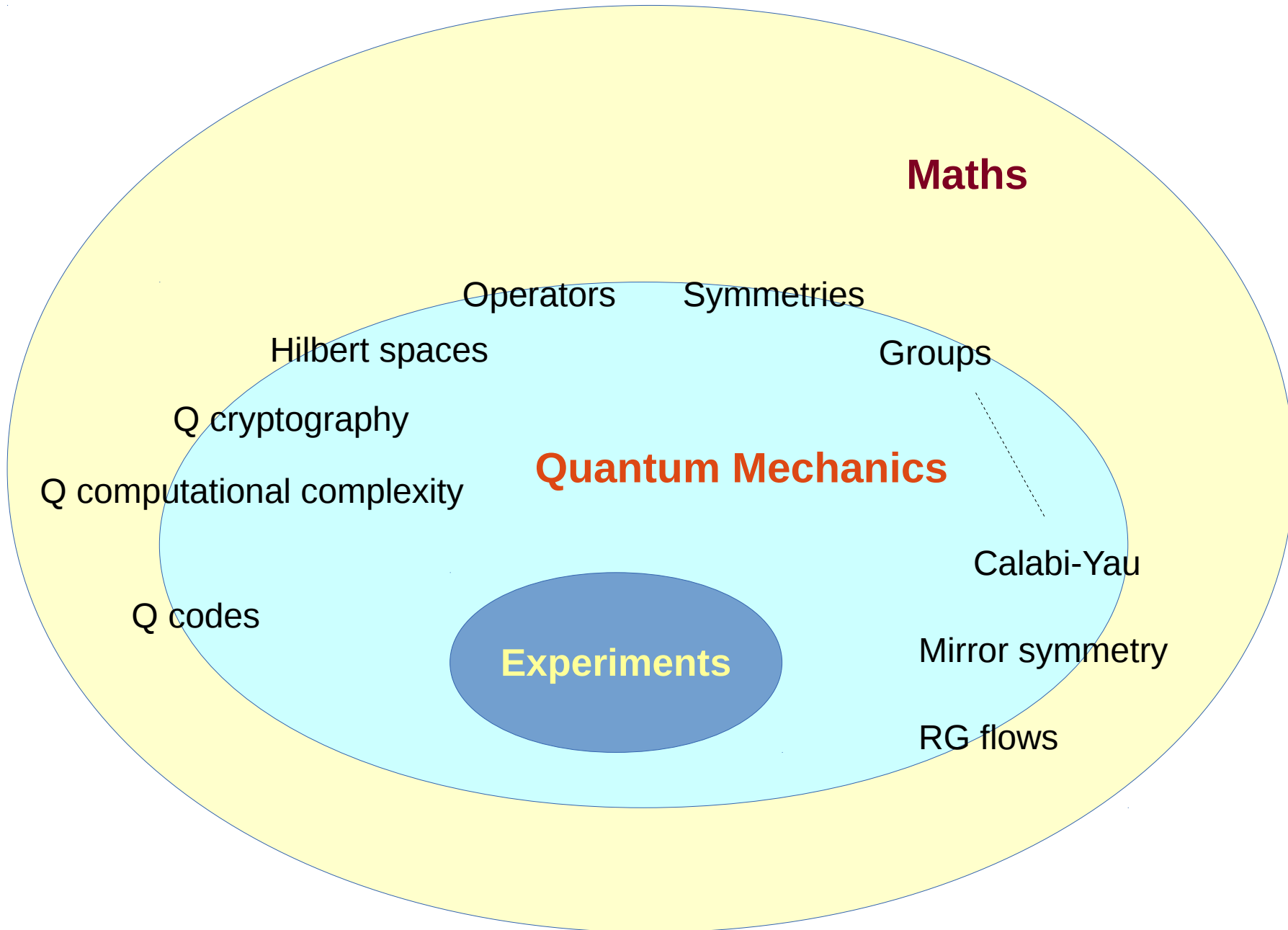
$i$	$j$	$i+j$	$i+2j$
0	0	0	0
0	1	1	2
0	2	2	1
1	0	1	1
1	1	2	0
1	2	0	2
2	0	2	2
2	1	0	1
2	2	1	0



# Divertimento on platonic entanglement



Tetrahedron	Faces = 4	AME(4,3)
Tetrahedron	Edges = 6	AME(6,5)
Tetrahedron	Vertices = 4	AME(4,3)
Exahedron	Faces = 6	AME(6,5)
Exahedron	Edges = 12	AME(12,11)
Exahedron	Vertices = 8	AME(8,7)
Octahedron	Faces = 8	AME(8,7)
Octahedron	Edges = 12	AME(12,11)
Octahedron	Vertices = 6	AME(6,5)
Dodecahedron	Faces = 12	AME(12,11)
Dodecahedron	Edges = 30	AME(30,29)
Dodecahedron	Vertices = 20	AME(20,19)
Icosahedron	Faces = 20	AME(20,19)
Icosahedron	Edges = 30	AME(30,29)
Icosahedron	Vertices = 12	AME(12,11)



# Why Quantum Mechanics?

Superposition  $|\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$

Entanglement  $|\psi\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$

Observables: Hermitian operators

$$O|\psi_\lambda\rangle = \lambda|\psi_\lambda\rangle$$

Measurements: Eigenvalues of Observables

Probability:  $|\langle\psi_\lambda|\psi\rangle|^2$

Evolution:  $|\psi(t)\rangle = e^{-\frac{i}{\hbar}(t-t_0)H}|\psi(t_0)\rangle$

# Falsifiability of Quantum Mechanics: Bell inequalities

Alice



Bob



$$P(\vec{A}, a; \vec{B}, b)$$

$$\vec{A} = \vec{n}_A \cdot \vec{\sigma}$$

$a$

measurement of spin in random directions

result of the measurement, take values +1 or -1

Measurements deliver a **probability distributions**

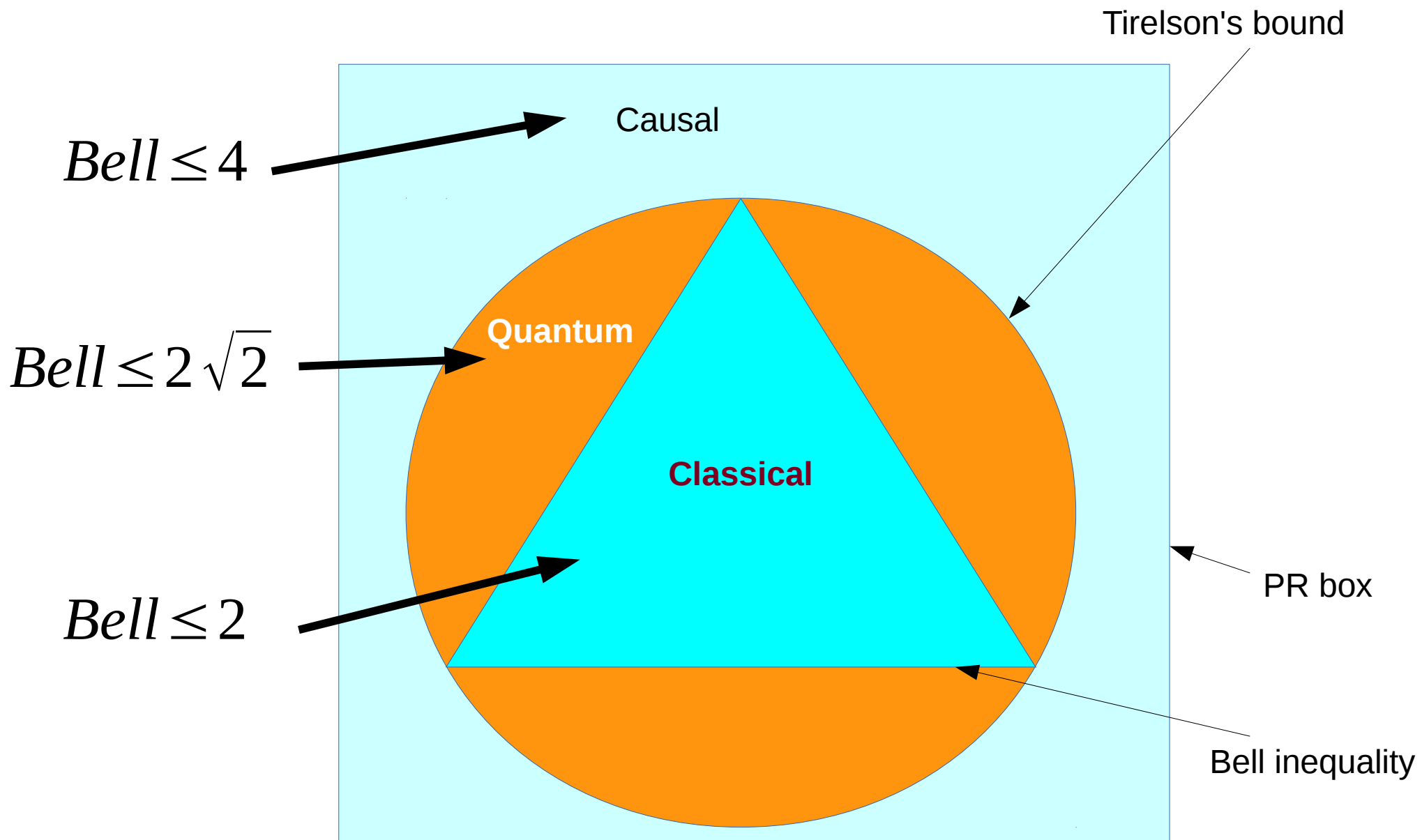
$$Bell = |\langle \vec{A}_1 \vec{B}_1 + \vec{A}_1 \vec{B}_2 + \vec{A}_2 \vec{B}_1 - \vec{A}_2 \vec{B}_2 \rangle|$$

In Classical Physics  $Bell \leq 2$

Einstein: element of reality, God does not play dice in the atoms



# Classical/Quantum/Causal polytopes for probabilities distributions



**NATURE: BELL < 2.82**

RG:  $H$  coarse grains in  $H'$

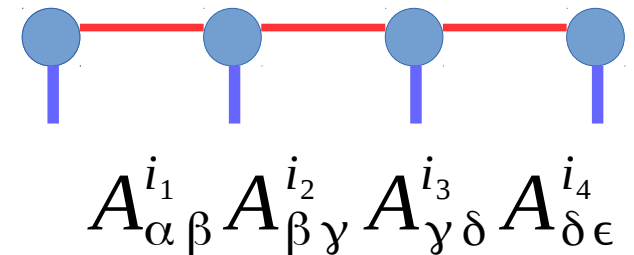
$H^*$  is a fixed point of the flow

RG on states

(Verstraete, Cirac, JIL, Rico, Wolf)

$$|\psi\rangle = \sum_{i_1, i_2, \dots, i_n} c^{i_1, i_2, \dots, i_n} |i_1, i_2, \dots, i_n\rangle$$

Matrix Product State (MPS)



$$A_{\alpha\beta}^{i_1} A_{\beta\gamma}^{i_2} = \sum_{\lambda} U_a^{i_1 i_2} \lambda^a V_{\alpha\gamma}^a$$

Coarse-graining

$$\lambda^a V_{\alpha\gamma}^a = \bar{A}_{\alpha\gamma}^{i_a}$$

## Goal of this talk

Produce some Quantum Information example of Holography

**Operators** in the **Bulk**  
dictated by  
operators on its **Boundary**

**State** in the **Bulk**  
dictated by  
its value on the **Boundary**